

# Bridging magic and non-Gaussian resources via Gottesman-Kitaev-Preskill encoding

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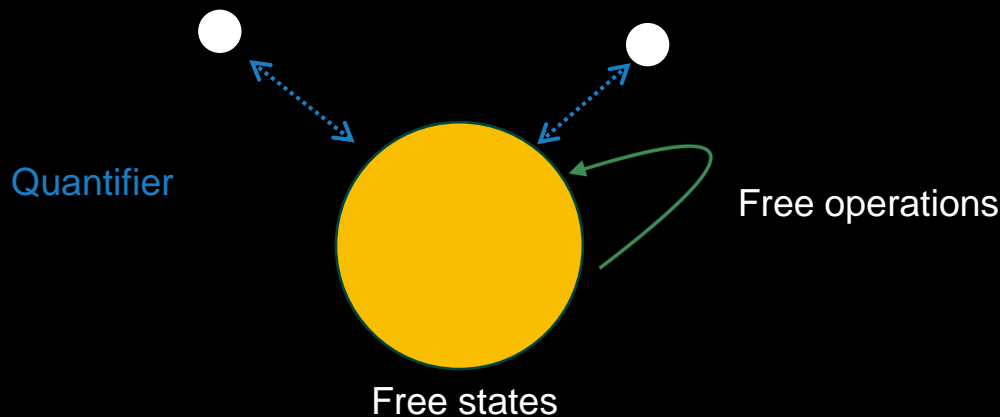
Joint work with Ryuji Takagi and Giulia Ferrini

arXiv:2406.06418

# Introduction

# Quantum Resources for quantum advantages

- **Big goal:** Quantitative understanding of quantum resources enabling quantum advantages underlying given physical and operational settings.
- **Quantum resource theories:** Framework to deal with **quantification** and **manipulation** of quantum resources

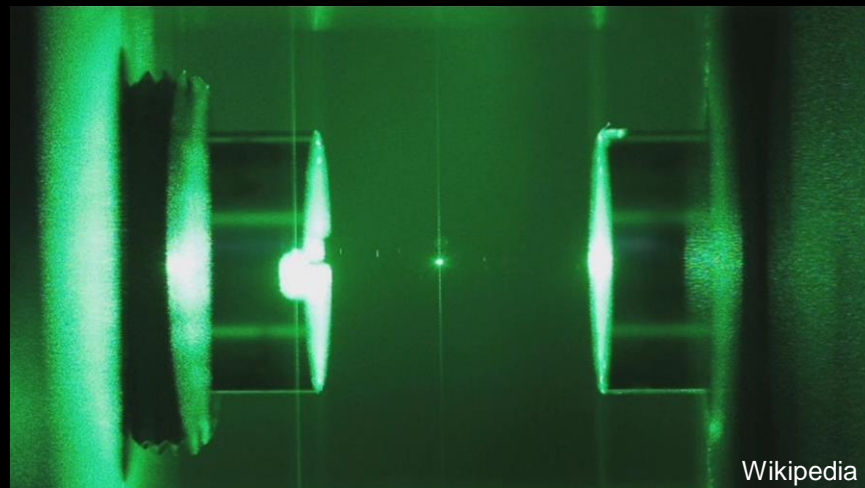


- Here, we study quantum computational resources relevant in CV and DV systems.

# Continuous Variables

- Quantum information encoded in q. modes
  - Harmonic oscillators
- Relevant observables ( $\hat{q}, \hat{p}$ ) have continuous spectrum
  - Infinite dimensional Hilbert space

$$[\hat{q}, \hat{p}] = i$$



# Gaussian quantum optics

- Displacement operators (CV Paulis):  $\hat{D}(\mathbf{r}) = \prod_{j=1}^n e^{-i r_{p_j} r_{q_j} / 2} e^{-i r_{q_j} \hat{p}_j} e^{i r_{p_j} \hat{q}_j}$

- Gaussian unitaries:  $\hat{U}_G \hat{D}(\mathbf{r}) \hat{U}_G^\dagger = \hat{D}(S\mathbf{r})$   $S\Omega S^T = \Omega$

- Important since it can be implemented experimentally
- Many analytical tools

$$\Omega = \begin{pmatrix} 0 & -\mathbb{1}_n \\ \mathbb{1}_n & 0 \end{pmatrix}$$

# Gaussian quantum optics

## Simulatability

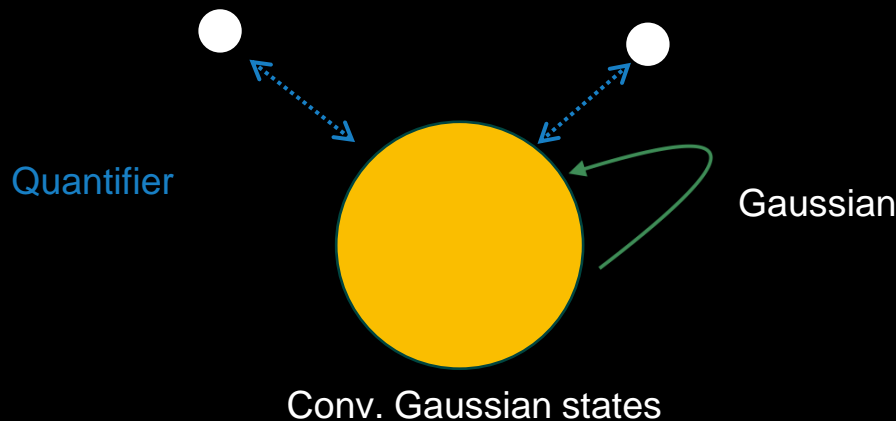
Any quantum process that begins with

- Gaussian states
  - Performs only Gaussian unitaries
  - Involves only measurements of canonical operators (including finite losses)
- can be simulated efficiently on a classical computer.

[Mari, Eisert, Phys. Rev. Lett. '12]

# Non-Gaussianity

- Gaussian circuits are classical simulatable
- Non-Gaussianity is a necessary resource for quantum advantage



[Takagi, Zhuang, Phys. Rev. A '18]

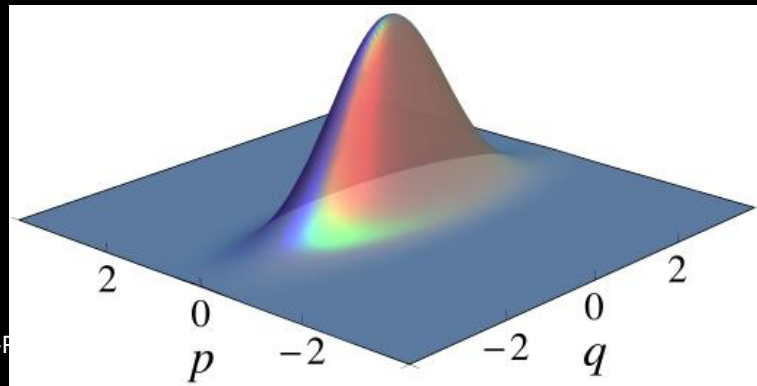
[Albarelli et al., Phys. Rev. A '18]

# CV Wigner function

- Phase-space representation of a quantum state
  - Fully equivalent to the density operator formalism

$$W_{\hat{\rho}}(\mathbf{r}) = \left(\frac{1}{2\pi}\right)^n \int_{\mathbb{R}^n} d\mathbf{x} e^{i\mathbf{r}_p \mathbf{x}} \left\langle \mathbf{r}_q + \frac{\mathbf{x}}{2} \left| \hat{\rho} \right| \mathbf{r}_q - \frac{\mathbf{x}}{2} \right\rangle_{\hat{q}}$$

- The Wigner function forms a quasi-probability distribution



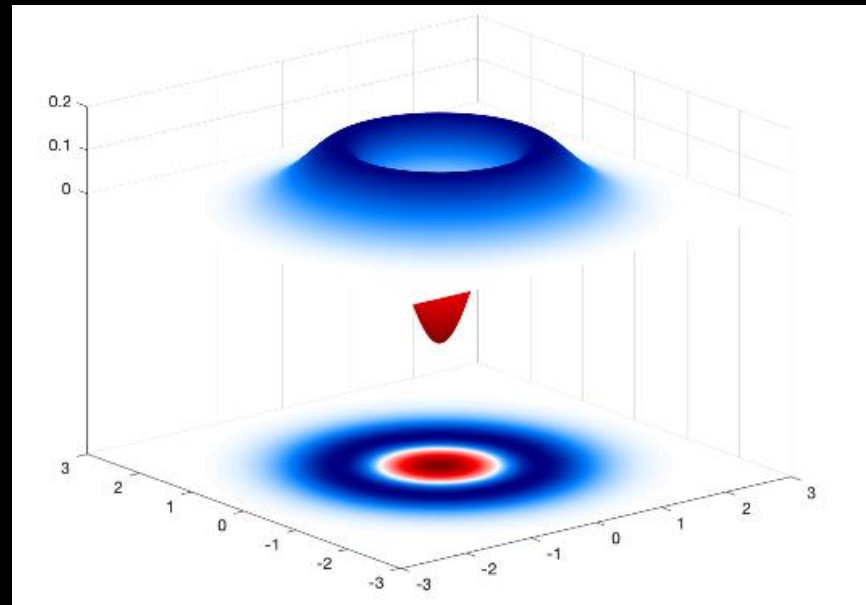


# Wigner negativity

- Wigner negativity shows genuine non-Gaussianity

$$\|W_{\rho}^{\text{CV}}\|_1 = \int d\mathbf{r} |W_{\rho}^{\text{CV}}(\mathbf{r})|$$

- Monotone under “Gaussian protocols”
  - Gaussian unitary
  - Attaching vacuum
  - Gaussian measurements
  - Gaussian feedforward



[Takagi, Zhuang, Phys. Rev. A '18]

[Albarelli et al., Phys. Rev. A '18]

# Discrete Variables

- Qudit Pauli  $\hat{P}_d(\mathbf{u}) = \bigotimes_{i=1}^n \omega_d^{\frac{1}{2}a_i b_i} \hat{X}_d^{a_i} \hat{Z}_d^{b_i}$

- A universal gate set is  $\{\hat{R}, \hat{P}, \underbrace{\hat{S}\hat{U}\hat{M}}_{\text{Clifford}}, \hat{T}\}$

- Clifford unitaries:  $\hat{U}_C \hat{P}_d(\mathbf{u}) \hat{U}_C^\dagger = \hat{P}_d(S\mathbf{u})$

# Why stabilizer and magic states?

## Gottesman-Knill theorem

A quantum computer based only on:

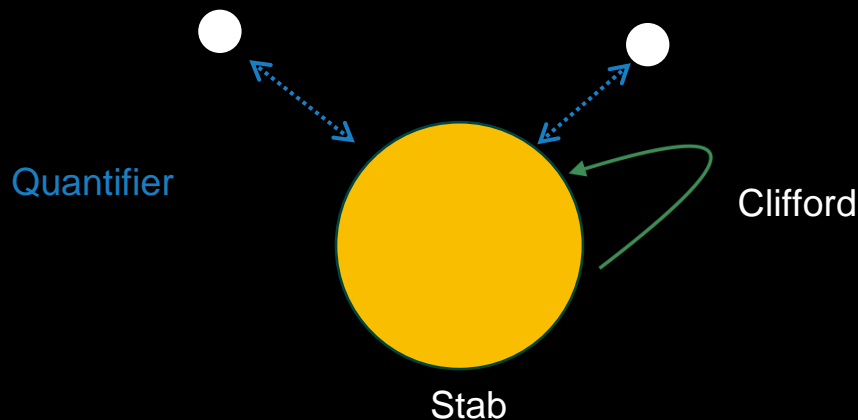
1. Qudits initialized in a Pauli eigenstate
2. Clifford group operations
3. Pauli measurements

Can be simulated efficiently with a classical computer

[Gottesman '98]

# Magic

- Stabilizer circuits are classical simulatable
- Magic is a necessary resource for quantum advantage



# DV Wigner function

## Odd Dimensions

- Phase-space representation of a DV quantum state
 
$$W_{\rho}^{\text{DV}}(\mathbf{u}) = d^{-n} \text{Tr} \left[ \hat{A}(\mathbf{u}) \hat{\rho} \right]$$

$$\hat{A}(\mathbf{u}) = d^{-n} \sum_{\mathbf{v} \in \mathbb{Z}_d^{2n}} \omega_d^{-\mathbf{u} \Omega_n \mathbf{v}^T} \hat{P}_d(\mathbf{v})^{\dagger}$$
- Wigner negativity  $\|W_{\rho}^{\text{DV}}\|_1 = \sum_{\mathbf{u}} |W_{\rho}^{\text{DV}}(\mathbf{u})|$
- Monotone under “Stabilizer protocols”
  - Clifford unitary
  - Auxilliary computation basis states
  - Pauli measurements
  - Clifford feedforward

[Veitch et al., New J. Phys. '14]

# Connecting CV and D

## DV

- Pauli
- Clifford
- DV Wigner function
- Magic
  - Negativity of Wigner function

## CV

- Displacements
- Gaussian
- CV Wigner function
- Non-Gaussianity
  - Negativity of Wigner function

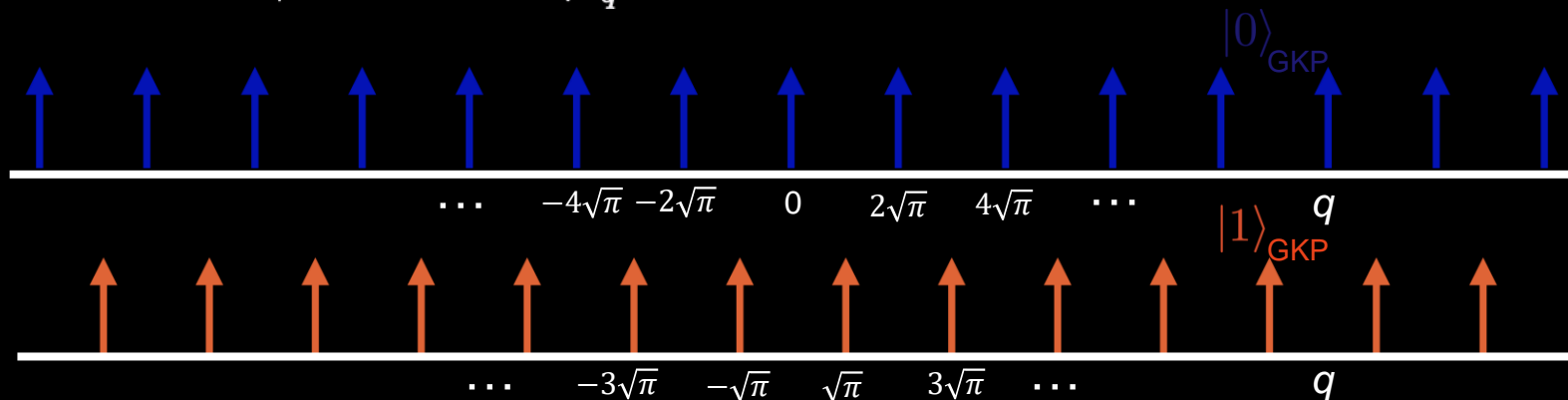
Any mapping between magic and non-Gaussian measures for a given state?

# Quantitative connections between DV and CV

# Gottesman-Kitaev-Preskill Code

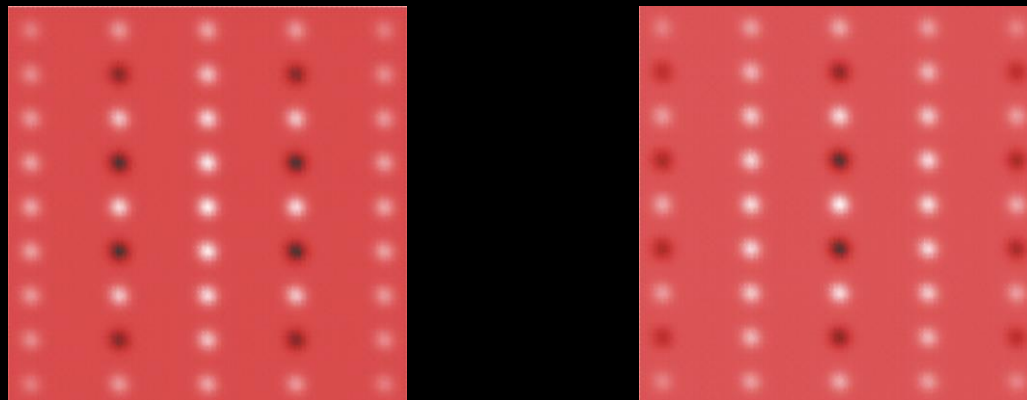
- Error correction code for continuous-variable systems that encode qudits

$$|j\rangle_{\text{GKP}} = \sum_{s=-\infty}^{\infty} \left| \sqrt{\frac{2\pi}{d}}(j + ds) \right\rangle_{\hat{q}}$$





# Wigner functions of GKP states

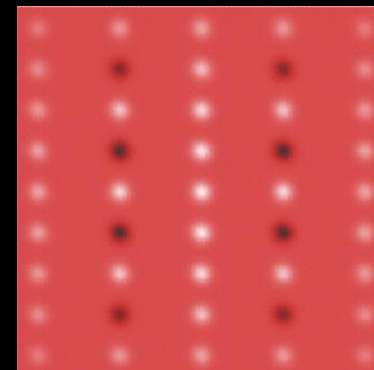
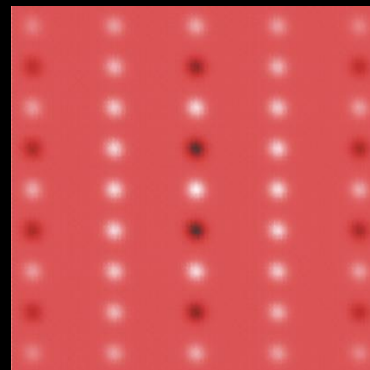


- (Ideal) GKP states are unnormalizable and have infinite non-Gaussianity
- The Wigner function is periodic with a unit cell  $[0, \sqrt{2d\pi})$

# Wigner functions of GKP states

- Consider negativity of one unit cell

$$\|W_{\rho_{\text{GKP}}}^{\text{CV}}\|_{1,\text{cell}} = \int_{\text{cell}} d\mathbf{r} |W_{\rho_{\text{GKP}}}^{\text{CV}}|$$



- Constant for pure stabilizer states
- Is there a connection between non-Gaussianity and magic?

[Yamasaki, Matsuura, Koashi, PRR '20]  
 [Hahn et al. PRL '22]

# CV-DV connection with Wigner function

- We consider the operator basis  $\hat{O}_{l,m} = \omega_d^{-ml/2} \hat{M}_l \hat{Z}_d^m$   $\hat{M}_l = \sum_{x \in \mathbb{Z}_d} |l-x\rangle \langle x|$
- Coefficients of this basis  $x_\rho(l, m) = d^{-n} \text{Tr}(\hat{O}_{l,m} \hat{\rho})$
- Wigner function

$$W_{\rho_{\text{GKP}}}^{\text{CV}}(\mathbf{r}) = \frac{\sqrt{d}^n}{\sqrt{8\pi}^n} \sum_{l,m} c_{\rho_{\text{GKP}}}(l, m) \delta\left(r_p - m\sqrt{\frac{\pi}{2d}}\right) \delta\left(r_q - l\sqrt{\frac{\pi}{2d}}\right)$$

$$\boxed{\underbrace{c_{\rho_{\text{GKP}}}(l, m)}_{\text{CV}} = \underbrace{x_\rho(l, m)}_{\text{DV}}}$$

# CV-DV connection with Wigner function

- General connection between Wigner negativity and magic
- For a n-qudit state we get

$$\|x_\rho\|_1 = \frac{\|W_{\rho_{\text{GKP}}}^{\text{CV}}\|_{1,\text{cell}}}{\|W_{\text{STAB}_n, \text{GKP}}^{\text{CV}}\|_{1,\text{cell}}}$$

Magic

Non-Gaussianity

- For odd:  $\|x_\rho\|_1 = \|W_\rho^{\text{DV}}\|_1$

# Magic measure

- For odd dimensions,  $\|x_\rho\|_1$  coincides with Wigner negativity  $\|W_\rho^{\text{DV}}\|$
- For even dimensions,  $\|x_\rho\|_1$  serves as a magic measure in the following sense:
  - Invariance under Clifford unitaries  $\|x_{U_C \rho U_C^\dagger}\|_1 = \|x_\rho\|_1$
  - Every pure stabilizer state  $\hat{\phi}$  takes the minimum value  $\|x_\phi\|_1 = 1$
  - For multi-qubit systems,  $\|x_\phi\|_1 = 1$  if and only if  $\hat{\phi}$  is a stabilizer state

# Application

- In GKP code, logical Clifford operations can be implemented by Gaussian operations.

$$|\bar{\psi}\rangle \xrightarrow{\text{H}} \bar{H} |\bar{\psi}\rangle \iff |\psi_{\text{GKP}}\rangle \xrightarrow{\text{G}} G |\psi_{\text{GKP}}\rangle$$

- Known implementations for logical non-Clifford operations use non-Gaussian operation

$$|\bar{\psi}\rangle \xrightarrow{\text{T}} \bar{T} |\bar{\psi}\rangle \iff \begin{matrix} |\psi_{\text{GKP}}\rangle \\ |\gamma\rangle \end{matrix} \xrightarrow{\text{G}} \bar{T} |\bar{\psi}\rangle$$

- Natural guess: Non-Gaussian operations are needed to implement logical non-Clifford gates.

Not obvious because GKP states already have non-Gaussianity initially.

# Non-Clifford needs non-Gaussianity

Suppose  $\Lambda$  is a channel with  $n$ -qubit input and output. If there is a pure magic state  $\hat{\psi}$  and a pure stabilizer state  $\hat{\phi}$  such that  $\Lambda(\hat{\phi}) = \hat{\psi}$ , then  $\Lambda$  cannot be implemented in the GKP code space by Gaussian protocols composed by

- Feedforwarded Gaussian operations conditioned on the measurement outcomes
- Gaussian unitary
- Attaching vacuum
- Gaussian measurements

- Extending previous finding for specific qubit operations
- For odd dimensions, the condition can be relaxed to the existence of a stabilizer

state  $\hat{\sigma}$  and a state  $\hat{\rho}$  such that  $\Lambda(\hat{\sigma}) = \hat{\rho}$  and  $\|W_{\rho}^{\text{DV}}\|_1 > 1$

[Yamasaki et al. Phys. Rev. Res. '20]

# Summary & Outlook

- Established the quantitative connection between magic (DV) and non-Gaussian (CV) by GKP encoding via Wigner and characteristic functions.
  - Showed that non-Clifford gate in GKP code space cannot be implemented by a Gaussian protocol
  - Proposed a simulation algorithm based on the distribution defined by a Hermitian extension of Pauli operators
- Finite squeezing?
- Other DV-CV connections with different bosonic codes?

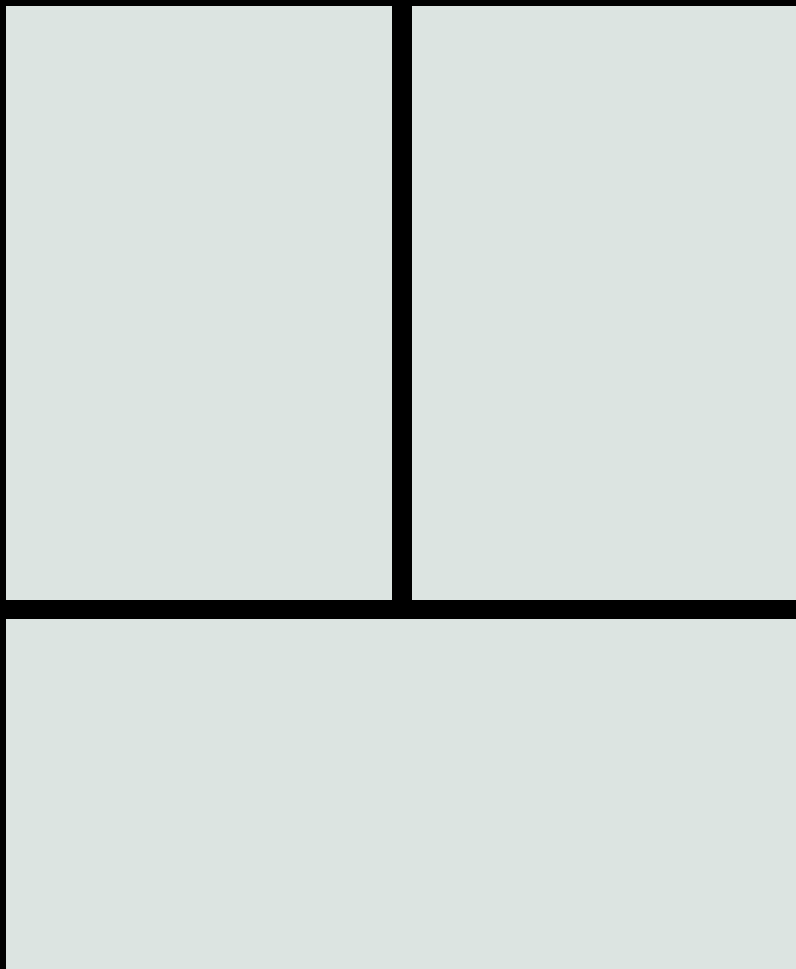


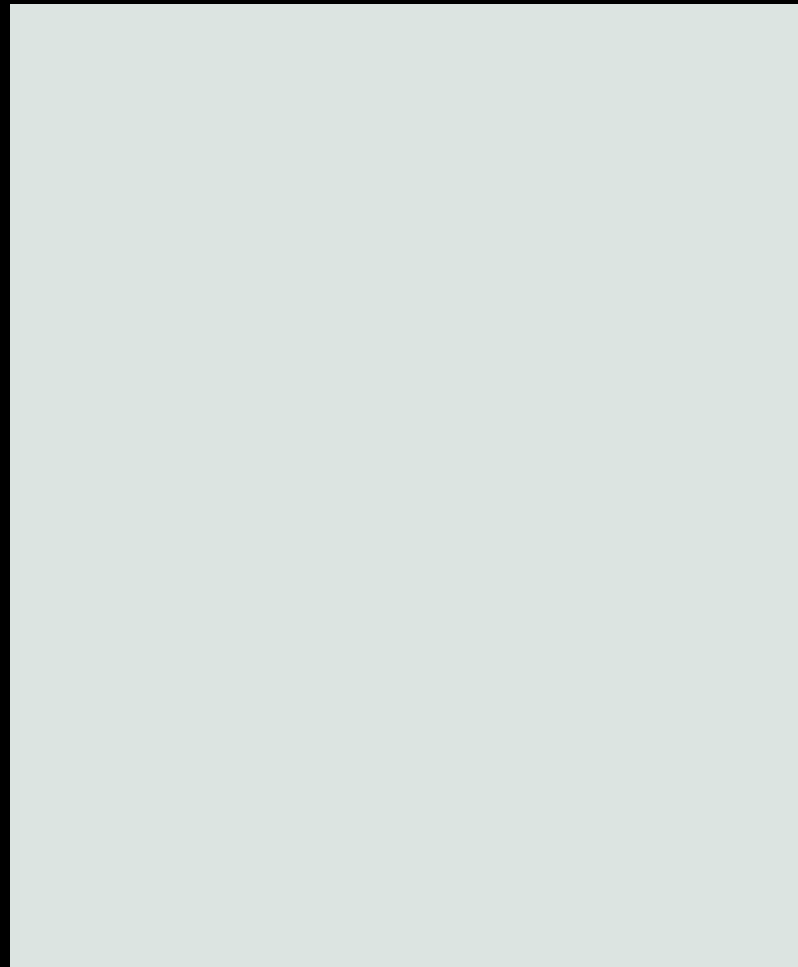
# Properties of $\mathcal{O}$

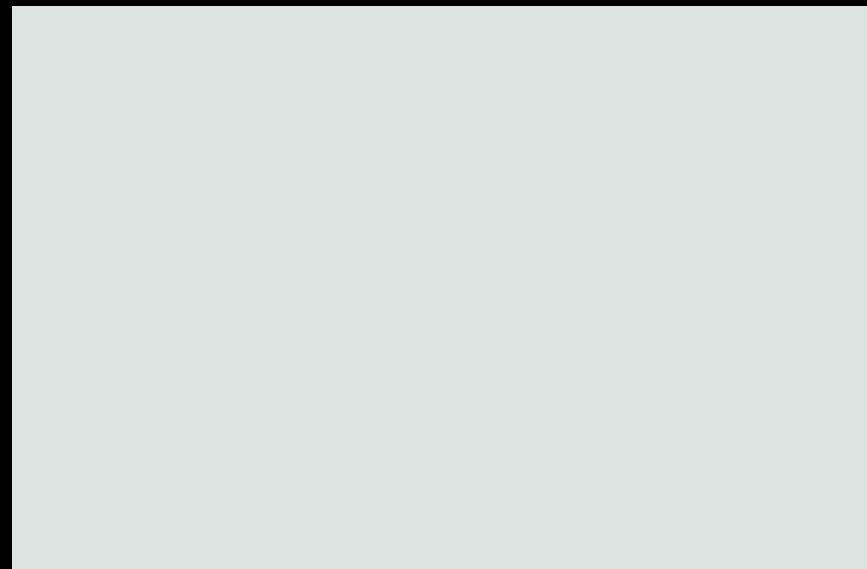
More results

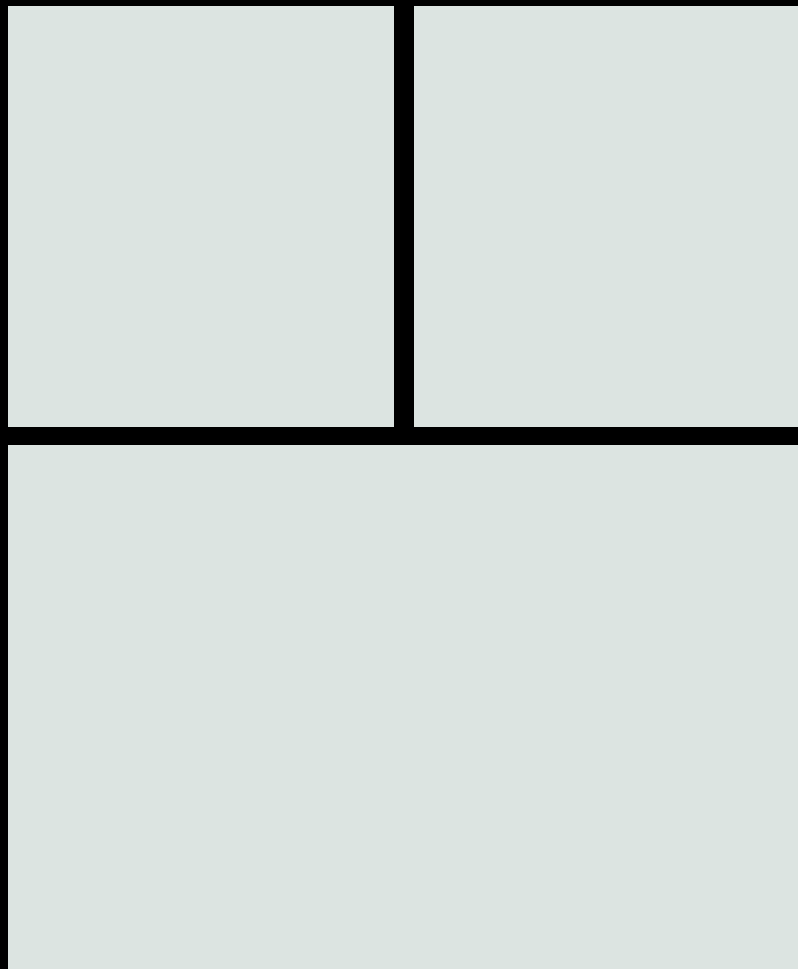
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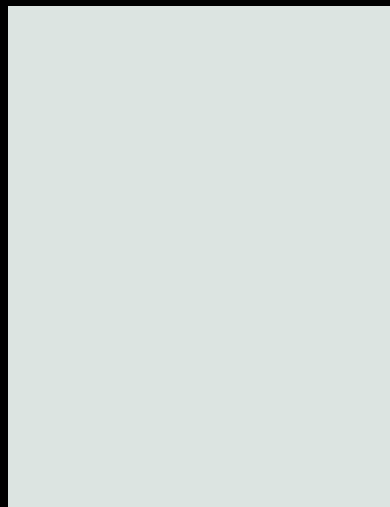
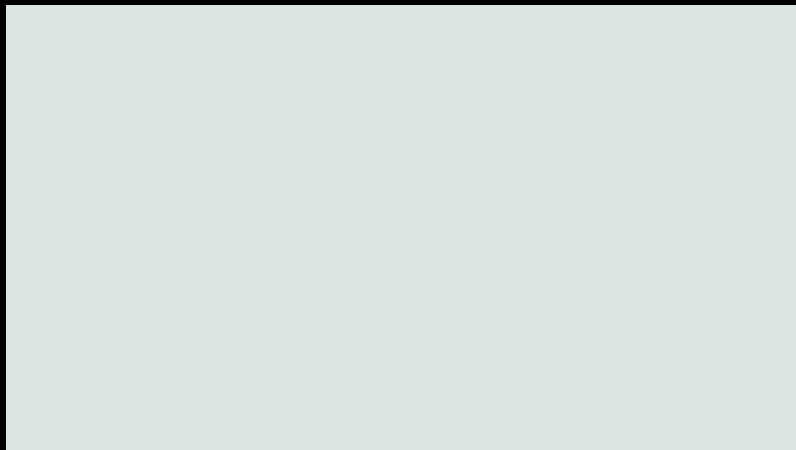




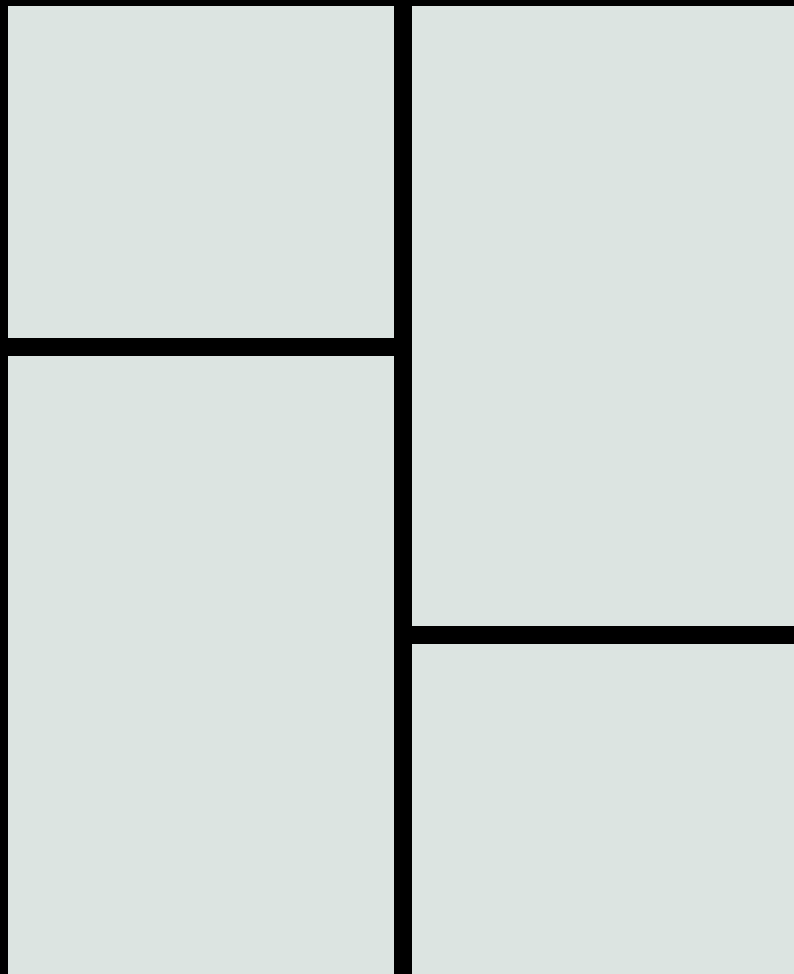










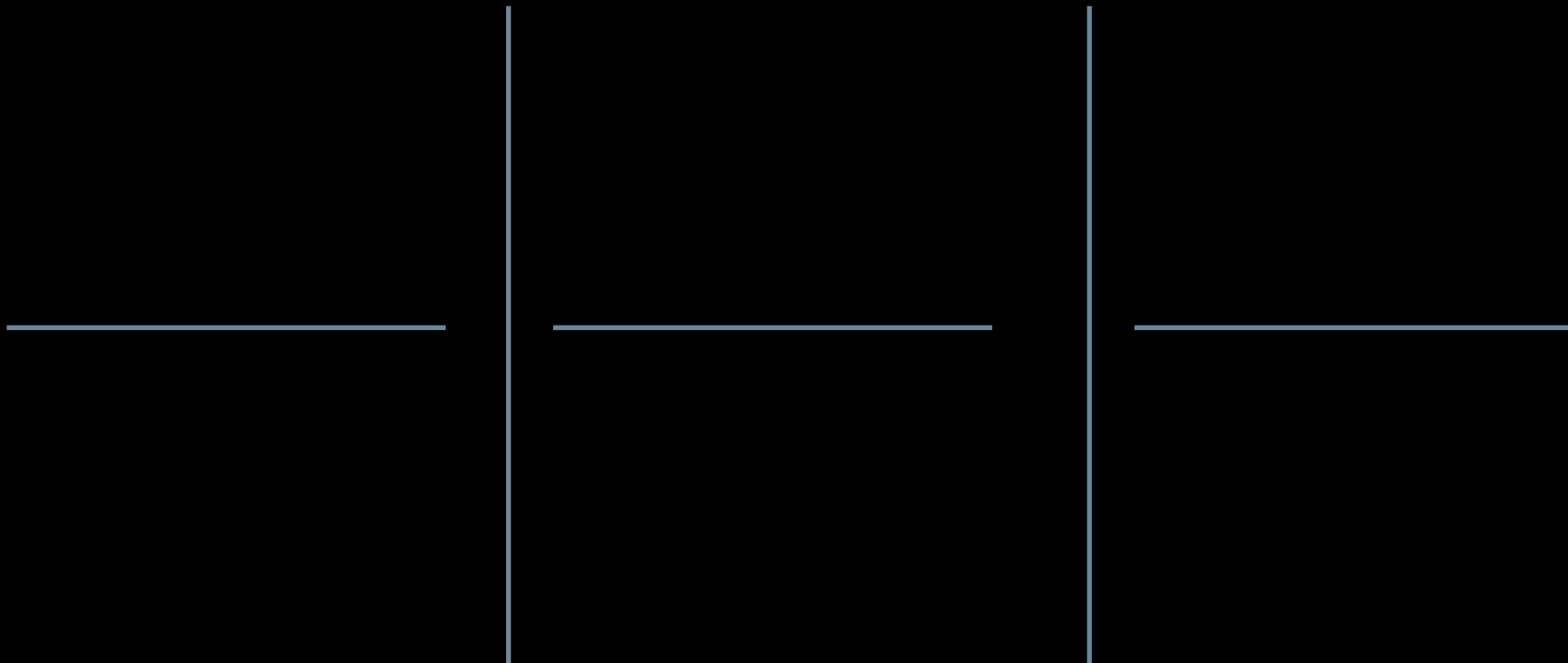




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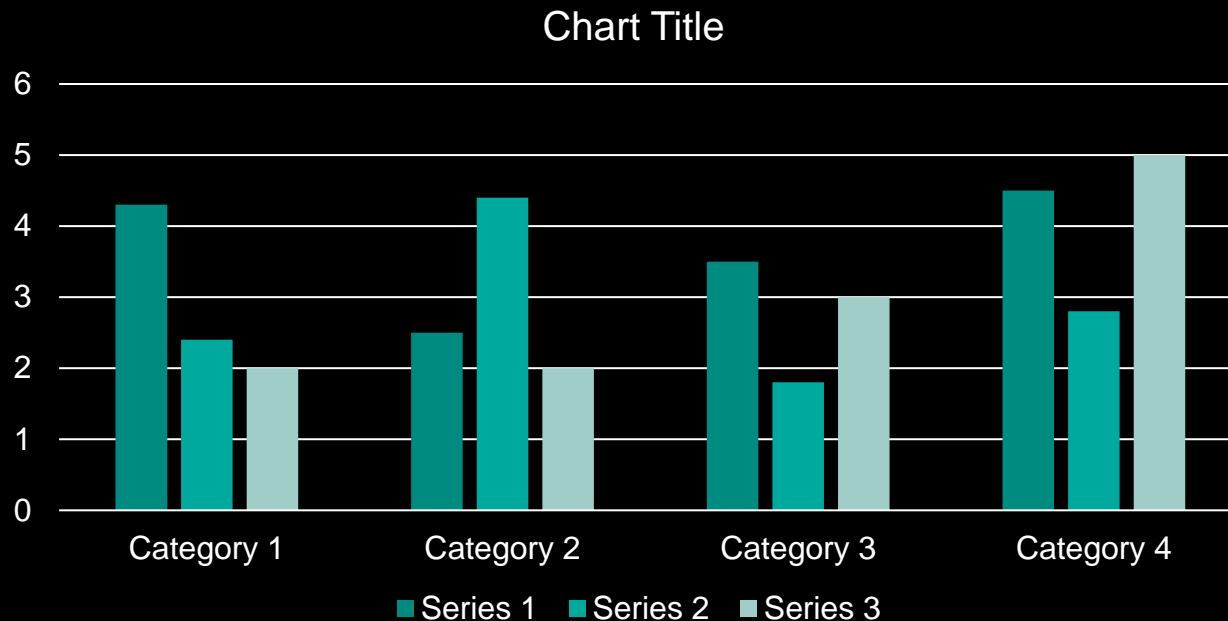
The following slides give examples of different charts. These are not intended to be used as templates but serve as guides and examples.

*Suggestion:*

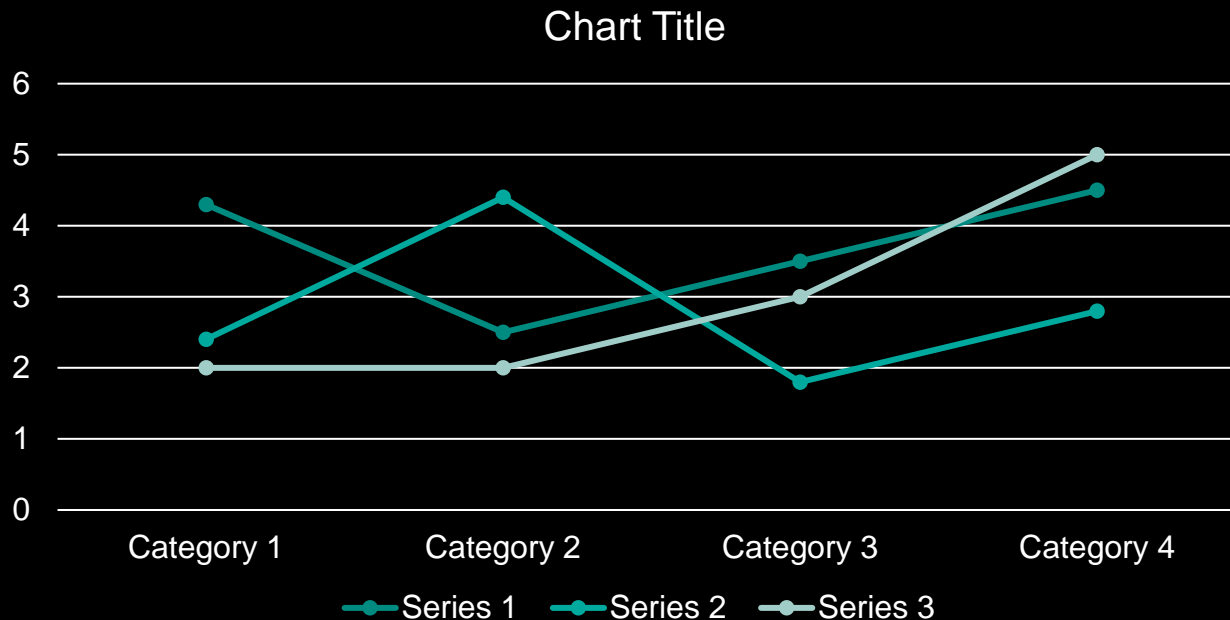
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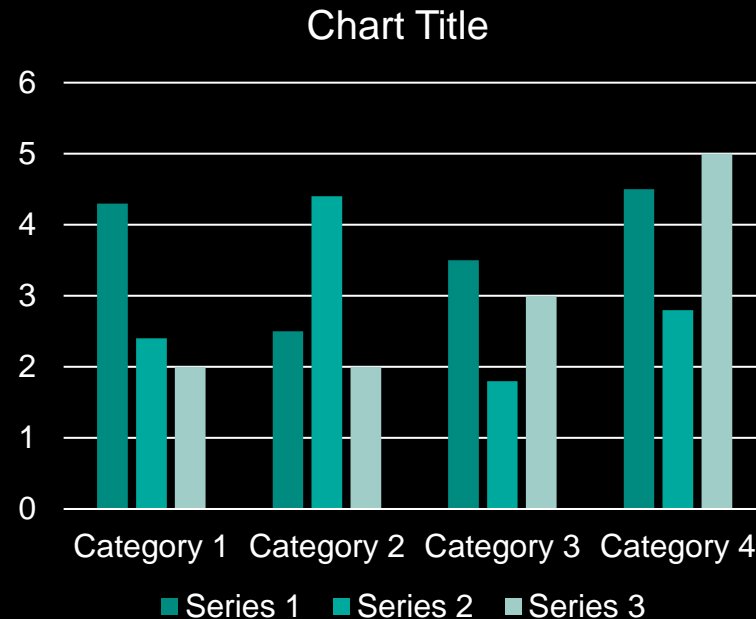
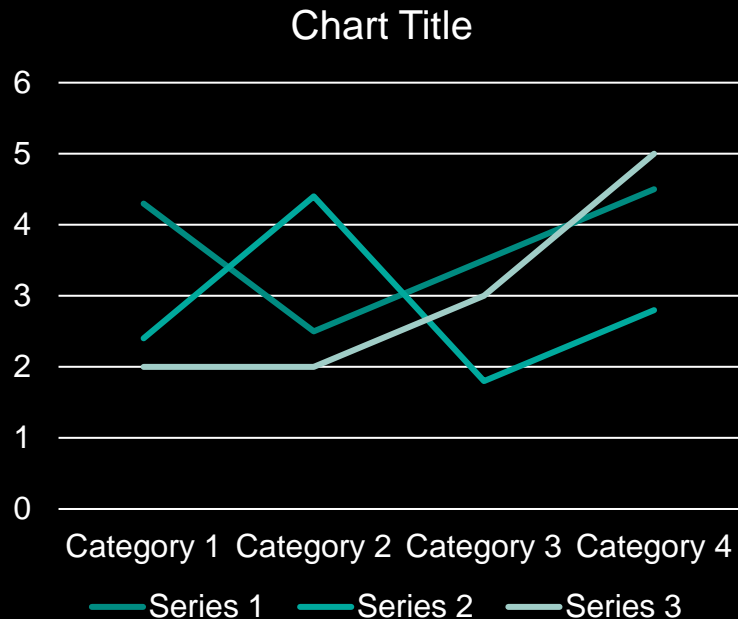
# Column chart example



# Line chart example

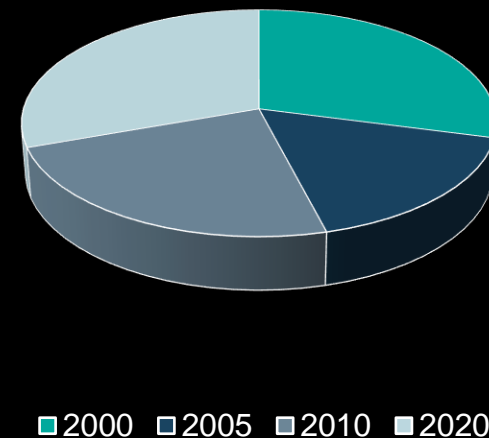
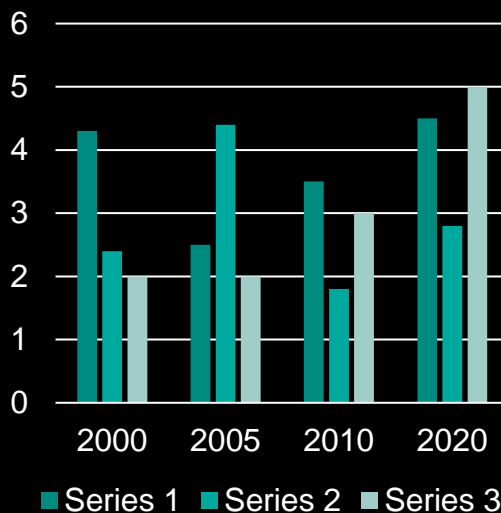
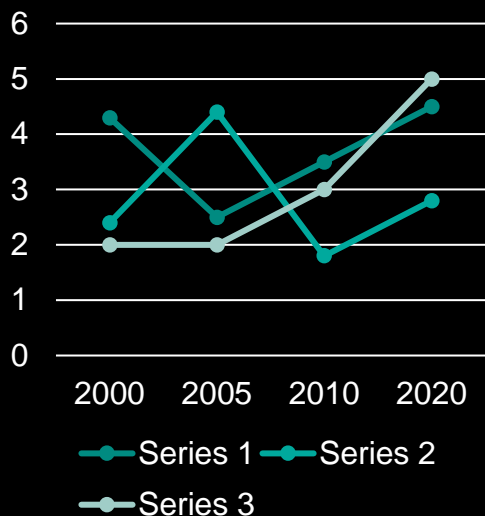


# Other chart type, two examples

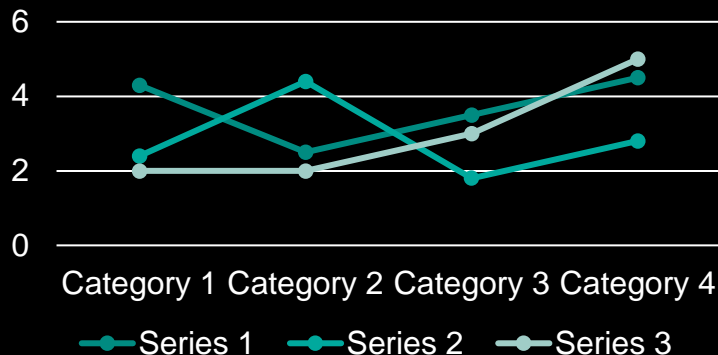




# Other chart type, three examples



# Chart and text examples



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# Table example

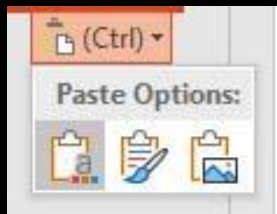
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